

QM 2

Oct 23

Today:

Interaction with EM field  
( $\sim 5.8.1 - 2$ )

- electric dipole interaction
- selection rules

WARM UP

$$\vec{E} = E_1 (\cos(\omega t), \sin(\omega t), 0)$$

$$\equiv E_0 \hat{\varepsilon} e^{-i\omega t} + E_0 \hat{\varepsilon}^* e^{i\omega t}$$

$$\text{let } \hat{\varepsilon}^* \cdot \hat{\varepsilon} = 1, E_0 \in \mathbb{R}$$

$$\text{Find } \hat{\varepsilon}, E_0$$

⇒ starting out with writing trig fns in LHS in  $e$ .

$$E_1 \left( \frac{e^{i\omega t} + e^{-i\omega t}}{2}, \frac{e^{i\omega t} - e^{-i\omega t}}{2i}, 0 \right)$$

$$\equiv E_0 \hat{\varepsilon} e^{-i\omega t} + E_0 \hat{\varepsilon}^* e^{i\omega t}$$

$$\Rightarrow \frac{E_1}{2i} \left[ i(e^{i\omega t} + e^{-i\omega t})\hat{x} + (e^{i\omega t} - e^{-i\omega t})\hat{y} + 0\hat{z} \right] = ( )$$

$$\Rightarrow \frac{E_1}{2i} \left[ e^{i\omega t} (i\hat{x} + \hat{y}) + e^{-i\omega t} (i\hat{x} - \hat{y}) \right] = ( )$$

comparing coefficients of  $e^{i\omega t}$  &  $e^{-i\omega t}$  on both sides:

$$\therefore \frac{E_1}{2i} (i\hat{x} + \hat{y}) = E_0 \hat{\epsilon}$$

$$\text{and, } \frac{E_1}{2i} (i\hat{x} - \hat{y}) = E_0 \hat{\epsilon}^*$$

dotting these equations to use the given condition  $\hat{\epsilon}^* \cdot \hat{\epsilon} = 1$

$$\Rightarrow -\frac{E_1^2}{4} (-1 - 1) = E_0^2$$

$$\Rightarrow \boxed{E_0 = \pm \frac{E_1}{\sqrt{2}}}$$

so  $\frac{E_1}{2i} (i\hat{x} + \hat{y}) = \pm \frac{E_1}{\sqrt{2}} \hat{z}$

normalising  $\Rightarrow \left\{ \begin{array}{l} \hat{z} = \mp i\sqrt{2} (i\hat{x} + \hat{y}) \\ \hat{z}^* = \pm i\sqrt{2} (-i\hat{x} + \hat{y}) \end{array} \right.$

$\left\{ \begin{array}{l} \hat{z} = \mp i \left( \frac{i\hat{x} + \hat{y}}{\sqrt{2}} \right) \\ \hat{z}^* = \pm i \left( \frac{-i\hat{x} + \hat{y}}{\sqrt{2}} \right) \end{array} \right.$

— x —

•  $\hat{z} = \frac{1}{\sqrt{2}} (\hat{e}_x + i\hat{e}_y)$  ,  $\hat{z}^* = \frac{1}{\sqrt{2}} (\hat{e}_x - i\hat{e}_y)$

→ these can be use as basis vectors for writing electric field

- \* spherical basis w.r.t Z axis
- \* useful for calculating angular momentum.

## Spherical basis

$$\hat{e}_1 = \frac{-1}{\sqrt{2}} (\hat{e}_x + i\hat{e}_y)$$

(minus sign would be convenient in calculations later)

$$\hat{e}_0 = \hat{e}_z$$

$$\hat{e}_{-1} = \frac{1}{\sqrt{2}} (\hat{e}_x - i\hat{e}_y)$$

$$\hat{e}_1^* = \frac{-1}{\sqrt{2}} (\hat{e}_x - i\hat{e}_y)$$

$$= -\hat{e}_{-1}$$

$$\hat{e}_0^* = \hat{e}_0$$

can be summarised with identity:

$$\boxed{\hat{e}_q^* = (-1)^q \hat{e}_{-q}}, \text{ with } q = 0, \pm 1$$

## electric dipole interaction (approximation)

- the interaction can be defined with some time-dependent perturbation:

$$V(t) \approx -q \vec{r} \cdot \vec{E}(\vec{r}=0, t)$$

- bound state,  $\vec{r} \approx 0$
- valid when  $\lambda \gg$  system size.
  - for atoms  $\lambda \sim 10^{-7} \text{ m}$ ,  $r \sim 10^{-10} \text{ m}$

(long wavelength approximation)

\* this is an approximation: we are looking at the electric field at only one point.

- dipole mom.:  $\vec{d} = q \vec{r}$   
 $\vec{d} = \sum_j q_j \vec{r}_j \rightarrow V = -\vec{d} \cdot \vec{E}(0, t)$

- apply perturbation:

- arbitrary polarization

$$\vec{E}(0,t) = E_0 \hat{\epsilon} e^{-i\omega t} + E_0^* \hat{\epsilon}^* e^{i\omega t}$$

- perturbation :

$$\begin{aligned} V(t) &= -\vec{d} \cdot \vec{E}(0,t) = e \vec{r} \cdot \vec{E}(0,t) \\ &= e E_0 (\vec{r} \cdot \hat{\epsilon}) e^{-i\omega t} + e E_0^* (\vec{r} \cdot \hat{\epsilon}^*) e^{i\omega t} \\ &= V^+ e^{-i\omega t} + V e^{i\omega t} \end{aligned}$$

$$\begin{aligned} \text{where, } V &= e E_0^* (\vec{r} \cdot \hat{\epsilon}^*) \\ V^+ &= e E_0 (\vec{r} \cdot \hat{\epsilon}) \end{aligned}$$

$$\text{consider, } \hat{\epsilon} = \hat{e}_q$$

→ possible transitions: - bound state → bound state

$$\begin{aligned} |i\rangle &= |n_i l_i m_i\rangle \\ \downarrow \\ |f\rangle &= |n_f l_f m_f\rangle \end{aligned}$$

(non-relativistic  
Hydrogen atom)

Transition Matrix Element:  $\langle f | V(t) | i \rangle$

- in harmonic pert. th., it was derived that:

$$\text{absorption: } \langle f | V^+ | i \rangle$$

$$\text{emission: } \langle f | V | i \rangle$$



absorption:

$$\begin{aligned} \langle f | V^+ | i \rangle &= \langle f | e E_0^* \vec{r} \cdot \hat{e}_q | i \rangle \\ &= e E_0^* \langle f | \vec{r} \cdot \hat{e}_q | i \rangle \end{aligned}$$

called the dipole matrix element.

$$\vec{r} \cdot \hat{e}_q = r_q = \begin{cases} -\frac{1}{\sqrt{2}}(x+iy), & q=1 \\ z, & q=0 \\ \frac{1}{\sqrt{2}}(x-iy), & q=-1 \end{cases}$$

minuses sign for  $q=1$  helps write  $r_q$  in terms of spherical harmonics.

components of  $\vec{r}$  in the spherical basis defined in WATK-UP.

$$= \sqrt{\frac{4\pi}{3}} Y_1^q(\theta, \phi) \dots \dots$$

$$\langle f | \vec{r} \cdot \hat{e}_q | i \rangle = \underbrace{\int dr r^2 r R_f(r) R_i(r)}_{\neq 0} \underbrace{\int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi Y_{l_f}^m(\theta, \phi)^* Y_1^q(\theta, \phi) Y_{l_i}^{m_i}(\theta, \phi) \sqrt{\frac{4\pi}{3}}}_{\equiv I \text{ Sakurai (3.393)}}$$

$$I \propto \langle l_f m_f | 1q, l_i m_i \rangle$$

$$= 0, \text{ unless } m_f = m_i + q$$

and  $l_f = |l_i - 1|, l_i, l_i + 1$   
and  $l_f$  &  $l_i$  have opposite parity

$$l_f = l_i \pm 1$$

Similarly for

Emission:

$$\langle f | \vec{r} \cdot \hat{e}_q^* | i \rangle = (-1)^q \langle f | \vec{r} \cdot \hat{e}_{-q}^* | i \rangle$$

$= 0$  unless:

$$\begin{cases} m_f = m_i - q \\ l_f = l_i \pm 1 \end{cases}$$

- for a general polarisation

$$\begin{aligned} m_f &= m_i + q, \quad q = 0, \pm 1 \\ l_f &= l_i \pm 1 \end{aligned}$$

$$\} \Delta n = \text{any}$$

Example: which are allowed:

a)  $|1\ 0\ 0\rangle \rightarrow |2\ 0\ 0\rangle$

b)  $|1\ 0\ 0\rangle \rightarrow |2\ 1\ 0\rangle$

c)  $|1\ 0\ 0\rangle \rightarrow |3\ 2\ 1\rangle$

⇒ a) not allowed (l cond<sup>n</sup> won't satisfy)

b) allowed

c) not allowed.



