

Nov-4
Today: Scattering (6.1, 6.2)

WARM UP: Scattering from $V(\vec{r})$

$$\frac{d\sigma}{d\Omega} \approx \left(\frac{mL^3}{2\pi\hbar^2} \right)^2 |V_{fi}|^2$$

$$\langle \vec{r} | \vec{k}_f \rangle = \frac{1}{L^{3/2}} e^{i\vec{k}_f \cdot \vec{r}}$$

$$\langle \vec{r} | \vec{k}_i \rangle = \frac{1}{L^{3/2}} e^{i\vec{k}_i \cdot \vec{r}}$$

a) write $\frac{d\sigma}{d\Omega}$ as integral on $V(\vec{r})$

b) for $V(\vec{r}) = g \delta(\vec{r})$, find $\frac{d\sigma}{d\Omega}$

$$\Rightarrow \text{a) } |V_{fi}|^2 = \left| \int \left(\frac{1}{L^{3/2}} e^{-i\vec{k}_f \cdot \vec{r}} \right) V \left(\frac{1}{L^{3/2}} e^{+i\vec{k}_i \cdot \vec{r}} \right) d\vec{r} \right|^2$$

$$\frac{d\sigma}{d\Omega} = \left(\frac{m_l^3}{2\pi\hbar^2}\right)^2 \left(\frac{1}{l^3}\right)^2 \left| \int e^{-i\vec{k}_f \cdot \vec{r}} N e^{i\vec{k}_i \cdot \vec{r}} d^3r \right|^2$$

(b)

$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \left(\frac{m}{2\pi\hbar^2}\right)^2 |g e^0|^2 \\ &= \left(\frac{m|g|}{2\pi\hbar^2}\right)^2 \end{aligned}$$

Quantum Wave Scattering:



$$\text{flux} = \frac{\text{prob.}}{\text{area} \times \text{time}} = \frac{1}{L^2 (L/V)} = \frac{\hbar k_i / m}{L^3}$$

then,

$$\text{detection rate} = \text{flux} \left(\frac{d\sigma}{d\Omega} \right) d\Omega$$

(from T.D.P.T)

- Outline (2 types of problems)

- 1) Born approximation.
 - 2) Partial wave expansion / scattering.
- formal Scattering theory

Zettili
develops
scattering
theory using
time independent
Schröd. eqn
instead of

* Derive from T.D.P.T (Sakurai)
* Transition Rate from $\vec{k}_i \rightarrow \vec{k}_f$

- $V(\vec{r})$ turned on slowly

$$\tilde{V}(\vec{r}, t) = e^{\varepsilon t} V(\vec{r}), \quad \varepsilon > 0$$

↳ potential (due to exponential form) is being turned on slowly from particle's perspective.

Time evolution operator:

$$\langle n | U_I(t, -\infty) | i \rangle = \delta_{ni} + \frac{1}{i\hbar} \int_{-\infty}^{\infty} dt' \langle n | V_I(t') U_I(t', -\infty) | i \rangle e^{\varepsilon t'}$$

T.D.P.T,
like Sakurai.

In the limit as $\varepsilon \rightarrow 0^+$

$$= \delta_{ni} + \frac{1}{i\hbar} \underbrace{T_{ni}}_{\text{some unknown (for now)}} \int_{-\infty}^t dt' e^{(i\omega_{ni} + \varepsilon)t'}$$

MOTIVATION: First order

$$\begin{aligned} \langle n | U_I(t, -\infty) | i \rangle &= \delta_{ni} + \frac{1}{i\hbar} \int_{-\infty}^t dt' \langle n | V_I(t') | i \rangle e^{\varepsilon t'} \\ &+ \delta_{ni} + \frac{1}{i\hbar} \int_{-\infty}^t dt' \langle n | e^{iH_0 t'/\hbar} V e^{-iH_0 t'/\hbar} | i \rangle e^{\varepsilon t'} \\ &= \delta_{ni} + \frac{1}{i\hbar} \int_{-\infty}^t dt' e^{i(\omega_{ni} + \varepsilon)t'} \underbrace{V_{ni}}_{T_{ni}} \end{aligned}$$

Similarly, 2nd order:

$$\langle n | U_I(t, -\infty) | i \rangle = \delta_{ni} + \frac{1}{i\hbar} \underbrace{\left(V_{ni} + \sum_m \frac{V_{nm} V_{mi}}{E_i - E_m + i\hbar\varepsilon} \right)}_{T_{ni}} \int_{-\infty}^t dt' e^{i(\omega_{ni} + \varepsilon)t'}$$

T-Matrix

$$T(E) = V + V \frac{1}{E - H_0 + i\hbar\varepsilon} T(E)$$

T-Matrix:

$$T_{ni}(E = E_i = E_n) = V_{ni} + \sum_m \frac{V_{nm} T_{mi}}{E_i - E_m + i\hbar\epsilon}$$

Scattering Cross Section

$$\frac{d\sigma}{d\Omega} = \left(\frac{mL^3}{2\pi\hbar^2} \right) |T_{fi}|^2 \quad (\text{exact})$$

—x—

Lippmann-Schwinger Equation

- Define "scattering state" $|\psi(t)\rangle$

$$T|i\rangle = V|\psi^{(+)}\rangle$$

Equation for $|\psi^{(+)}\rangle$

$$T = V + V G_0 T$$

$$G_0 = \frac{1}{E - H_0 + i\hbar\epsilon}$$

$$T|i\rangle = V|i\rangle + V G_0 T|i\rangle$$

$$V|\psi^{(+)}\rangle = V|i\rangle + V G_0 V |\psi^{(+)}\rangle$$

$$|\psi^{(+)}\rangle = |i\rangle + G_0(E) V |\psi^{(+)}\rangle$$

$$|\psi^{(+)}\rangle = |i\rangle + \frac{1}{E - H_0} V |\psi^{(+)}\rangle \quad \eta \rightarrow 0^+$$

Lippmann
Schwinger
eqn

$|\psi^{(+)}\rangle$ solves the T.I.S.E:

$$(E - H_0 + i\eta) |\psi^{(+)}\rangle = (E - H_0 + i\eta) |i\rangle + V |\psi^{(+)}\rangle$$

$\swarrow \quad \swarrow$
 $E_i - E_i = 0 \quad \rightarrow \lim_{\eta \rightarrow 0^+}$
 $\searrow$
 ≈ 0

$$\therefore (E - H_0) |\psi^{(+)}\rangle = V |\psi^{(+)}\rangle$$

$$E |\psi^{(+)}\rangle = (H_0 + V) |\psi^{(+)}\rangle$$

L.S.E in Integral form

$$\psi^{(+)}(\vec{r}) = \langle \vec{r} | \psi^{(+)} \rangle = \langle \vec{r} | i \rangle +$$

$$\underbrace{\phi_i(\vec{r})}$$

$$\langle \vec{r} | \frac{1}{E_i - H_0 + i\eta} V | \psi^{(4)} \rangle$$

↑ insert identity operators.

$$\mathbb{I} = \int d^3\vec{r}' |\vec{r}'\rangle \langle \vec{r}'|$$

$$= \phi_i(\vec{r}) + \int d^3\vec{r}' \langle \vec{r} | \frac{1}{E_i - H_0 + i\eta} | \vec{r}' \rangle V(\vec{r}') \psi^{(4)}(\vec{r}')$$

Green's
function

Define :

$$G_+(\vec{r}, \vec{r}') \equiv \frac{\hbar^2}{2m} \langle \vec{r} | \frac{1}{E_i - H_0 + i\eta} | \vec{r}' \rangle$$

$$\Rightarrow \psi^{(4)}(\vec{r}) = \phi_i(\vec{r}) + \frac{2m}{\hbar^2} \int d^3\vec{r}' G_+(\vec{r}, \vec{r}') V(\vec{r}') \psi^{(4)}(\vec{r}')$$

Lindgren L.S.E





