

Nov 11

Scattering Amplitude (6.2)

WARM UP: Periodic B.C. Box

$$\langle \vec{r} | \vec{k} \rangle = \frac{1}{L^{3/2}} e^{i \vec{k} \cdot \vec{r}}$$

$$\vec{k} = \frac{2\pi}{L} (n_x, n_y, n_z)$$

Find

$$\int d^3 \vec{k} f(\vec{k}) \approx (?) \sum_{\vec{k}} f(\vec{k})$$

HINT

→

$$\sum_n \Delta x f(x) \approx \int dx f(x)$$

have to find
this

as $\Delta x \rightarrow 0$

$$\Delta k = \frac{2\pi}{L}$$

$$(\Delta k)^3 = \left(\frac{2\pi}{L}\right)^3 \quad (\because \text{the integral is 3-D})$$

$$\Rightarrow \quad ? \rightarrow \left(\frac{2\pi}{L}\right)^3$$

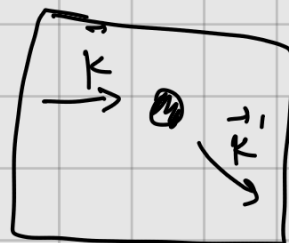
$$\therefore \int d^3k f(k) = \left(\frac{2\pi}{L}\right)^3 \sum_k f(k)$$

Review:

$$\text{initial } |i\rangle : \langle \vec{r} | i \rangle = \frac{1}{L^{3/2}} e^{i\vec{k} \cdot \vec{r}}, \quad E = \frac{\hbar^2 k^2}{2m}$$

Potential : $V(\vec{r})$

$$\text{final } |f\rangle : \langle \vec{r} | f \rangle = \frac{1}{L^{3/2}} e^{i\vec{k}' \cdot \vec{r}}$$



* we don't know what $\vec{k}'$ is (b/c Quantum Mech!), but we can predict its probability distribution.

Differential Cross-section

$$\frac{d\sigma}{d\Omega} = \left(\frac{mL^3}{2\pi\hbar^2}\right)^2 |T_{fi}|^2$$

$$\text{where } : T = V + V G_0 T$$

describes
propagation
(how?)
it's called
resolvent in
this form

$$G_0(E) = \frac{1}{E - H_0 + i\eta} \quad (\eta \rightarrow 0^+)$$

$$H_0 = \frac{\vec{p}^2}{2m} = -\frac{\hbar^2}{2m} \nabla^2$$

↑
(operator)

Scattering states

$$T_{fi} = \langle f | T | i \rangle = \langle f | V | \psi^{(+)} \rangle$$

Find $|\psi^{(+)}\rangle$, satisfies

- $(H_0 + V) |\psi^{(+)}\rangle = E |\psi^{(+)}\rangle, E \geq 0$

- L.S.E

$$\psi^{(+)}(\vec{r}) = \frac{1}{L^{3/2}} e^{i\vec{k} \cdot \vec{r}} + \frac{2m}{\hbar^2} \int d^3\vec{r}' G_+(\vec{r}, \vec{r}') V(\vec{r}') \psi^{(+)}(\vec{r}')$$

$$G_+(\vec{r}, \vec{r}') = \frac{\hbar^2}{2m} \langle \vec{r} | \frac{1}{E - H_0 + i\eta} | \vec{r}' \rangle$$

$$\mathbb{I} = \sum_{\vec{k}'} |\vec{k}'\rangle \langle \vec{k}'|$$

Green's function

$$G_+(\vec{r}, \vec{r}') = \frac{\hbar^2}{2m} \langle \vec{r} | \frac{1}{E - H_0 + i\eta} \sum_{\vec{k}'} |\vec{k}'\rangle \langle \vec{k}'| | \vec{r}' \rangle$$

projector.

$$= \frac{\hbar^2}{2m} \sum_{\vec{k}'} \langle \vec{r} | \frac{1}{\left(\frac{\hbar^2 k^2}{2m}\right) - \left(\frac{\hbar^2 (k')^2}{2m}\right) + i\eta} | \vec{k}' \rangle \frac{1}{L^3} e^{i\vec{k}' \cdot \vec{r}}$$

$$= \sum_{\vec{k}'} \underbrace{\frac{1}{L^3}}_{\substack{\downarrow \\ \frac{1}{(2\pi)^3} \text{ (from warm-up)}}} e^{i\vec{k}' \cdot \vec{r}} \frac{1}{k^2 - (k')^2 + i\varepsilon} e^{-i\vec{k}' \cdot \vec{r}'}$$

$$= \frac{1}{(2\pi)^3} \int d^3 \vec{k}' \frac{e^{i\vec{k}' \cdot (\vec{r} - \vec{r}')}}{k^2 - (k')^2 + i\varepsilon}$$

$$G_+(\vec{r}, \vec{r}') = \int \frac{d^3 k'}{(2\pi)^3} \frac{e^{i\vec{k}' \cdot (\vec{r} - \vec{r}')}}{k^2 - (k')^2 + i\varepsilon}$$

$$k^2 - (k')^2 + i\varepsilon = -(k' - a)(k' - b) \quad (\text{factoring})$$

$$= -(k')^2 + (b+a)k' - ab$$

by comparison

$$\Rightarrow a+b=0, \quad ab=i\varepsilon$$

$$\underbrace{\hspace{10em}}_{a=\sqrt{k^2+i\varepsilon}, \quad b=-\sqrt{k^2+i\varepsilon}}$$

$$\Rightarrow a = k \sqrt{1+i\frac{\varepsilon}{k^2}}, \quad b = -k \sqrt{1+i\frac{\varepsilon}{k^2}}$$

Since ϵ is small,
Taylor expanding the R.H.S

$$\rightarrow \boxed{a \approx k + \frac{i\epsilon}{2k^2}}$$

$$a = k + i\epsilon$$

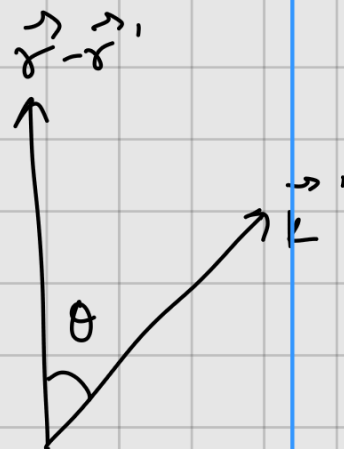
$$\boxed{b \approx -k - \frac{i\epsilon}{2k^2}}$$

$$a = -k - i\epsilon \quad (\text{say})$$

$$\therefore G_+(\vec{r}, \vec{r}') = - \int \frac{d^3 k'}{(2\pi)^3} \frac{e^{i\vec{k}' \cdot (\vec{r} - \vec{r}')}}{(k' - (k + i\epsilon))(k' + (k - i\epsilon))}$$

*

let z axis be along $\vec{r} - \vec{r}'$



$$\Downarrow \int_0^\infty \frac{(k')^2 dk'}{(2\pi)^3} \int_0^\pi \frac{d\theta \sin\theta 2\pi e^{i\vec{k}' \cdot (\vec{r} - \vec{r}') \cos\theta}}{(k' - (k + i\epsilon))(k' + (k - i\epsilon))} \quad (-1)$$

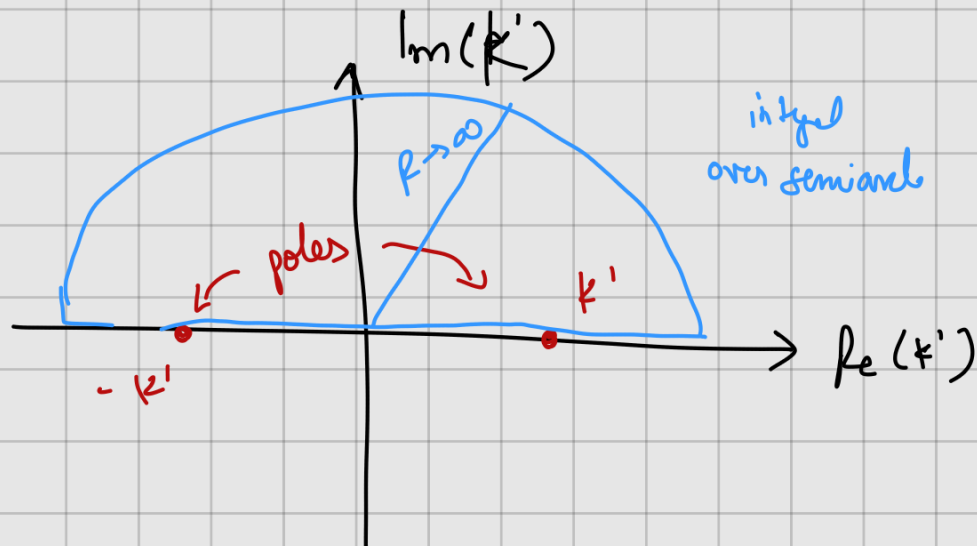
$$= \int \frac{(k')^2 dk'}{(2\pi)^3} 2\pi \int_{-1}^1 du \frac{e^{ik' |\vec{r} - \vec{r}'| u}}{[k' - a][k' + a]}$$

$$= \int_0^{\infty} \frac{(k')^2 dk'}{(2\pi)^2} \frac{e^{ik'|\vec{r}-\vec{r}'|} - e^{-ik'|\vec{r}-\vec{r}'|}}{(ik')(|\vec{r}-\vec{r}'|)(k'-a)(k'+a)} \quad (-1)$$

$$G_+(\vec{r}, \vec{r}') = \frac{i}{|\vec{r}-\vec{r}'|} \int_0^{\infty} \frac{k' dk'}{(2\pi)^2} \frac{e^{ik'|\vec{r}-\vec{r}'|}}{(k'-(k+i\delta))(k'+(k-i\delta))}$$

Using Cauchy residue theorem

$$\oint \frac{f(z) dz}{z-z_0} = 2\pi i f(z_0) \quad \hookrightarrow k'$$



$$G_+(\vec{r}, \vec{r}') = \frac{i}{|\vec{r}-\vec{r}'|} \oint \frac{z dz}{(2\pi)^2} \frac{e^{iz|\vec{r}-\vec{r}'|}}{(z-(k+i\delta))(z-(-k-i\delta))}$$

$$= \frac{i}{|\vec{r}-\vec{r}'|} (2\pi i) \frac{(k+i\delta)}{(2\pi)^2} \frac{e^{i(k+i\delta)|\vec{r}-\vec{r}'|}}{(k+i\delta)-(-k-i\delta)}$$

$$G_+(\vec{r}, \vec{r}') = - \frac{1}{4\pi |\vec{r} - \vec{r}'|} e^{ik |\vec{r} - \vec{r}'|}$$

